

Growth of Clathrate Hydrates from Water Drops in Cyclopentane

Jorge Peixinho,*¹ Valentin Ageorges, and Benoit Duchemin

Laboratoire Ondes et Milieux Complexes, CNRS et Université Le Havre Normandie, 76600 Le Havre, France

ABSTRACT: Clathrate hydrates are icelike crystalline compounds with encaged guest molecules trapped inside the cages of hydrogen-bonded water molecules. Their growth is visualized for cyclopentane within water drops of 2 μL on glass and polytetrafluoroethylene surfaces. The effect of the interfacial tension between water and cyclopentane is measured at different temperatures and for different concentrations of an oil-soluble surfactant, sorbitan monooleate (Span 80). The drops experience a temperature sequence where they freeze into ice, form hydrates, and melt. Cyclopentane hydrate crystals are affected by the concentration of the surfactant. The morphology seen here could be relevant for explaining the behavior of hydrate emulsions.

INTRODUCTION

Clathrate hydrates¹ (abbreviated hydrates) are crystalline solids composed of water molecules called the “host” that form cages in which “guest” molecules are enclosed. The guests may be light hydrocarbons, cyclopentane (CP), carbon dioxide, methane, rare gas, etc. It is generally acknowledged that hydrates have properties that include a large capacity of gas storage, fractionation of gas mixtures, and high heats of formation and decomposition. These properties enable hydrates to be used for various technologies in transportation and storage of natural gas, sequestration of carbon dioxide, and even water desalination.² There are several types of hydrates that differ in crystallographic cage structure, types I, II, and H, in which nature and size depend upon the guest molecule. CP guests¹ often lead to the formation of type II.

Most previous works were concerned with the kinetics of hydrate crystallization under different pressure and temperature conditions.³ Typically, hydrates are formed at temperatures below about 4 °C (277 K) in combination with elevated pressure in the range of 1–10 MPa. However, not all hydrate-forming systems require elevated pressures, and both CP and tetrahydrofuran (THF) form hydrates at atmospheric pressure and accessible temperatures. For example, the melting temperature of CP hydrate is approximately 7 °C (280 K) at atmospheric pressure.⁴ This allows for the study of hydrate formation processes without the experimental difficulties of dealing with pressure. Because THF is nearly fully miscible with water, the essential issue of mass transfer of the hydrate former (guest species) from the external organic phase to the water drops is not present in a THF–water system, whereas CP–water systems retain an interface.

The growth of hydrates deserves to be studied because of its implication in the blockage of pipelines, i.e., flow assurance. Previous investigations on hydrate suspensions in oil-dominated systems,⁵ water-in-crude oil emulsions,⁶ and black oil suspensions⁷ have tested the effects of the temperature and conversion rate of water into hydrates. However, the detailed effect of the component of each system is sometimes difficult to measure quantitatively. Hence, model systems of water-in-oil emulsions are used to mimic the hydrate growth. Specifically, the presence of hydrates is detected from the significant increase of the viscosity seen in rheological measurements

during temperature quench. This overshoot cannot be explained by the simple conversion of water drops into ice or hydrates and has been used to explore the effect of different parameters controlling the formation of hydrates for different subcooling temperatures,^{8,9} shear rates,^{9–11} water fractions,^{8,9,11–13} surfactant concentrations,⁹ cooling rates,¹² and salt concentrations.¹¹ A mechanism based on random nucleation of water droplets into a porous hydrate particle,¹⁴ agglomeration,¹⁵ and capillary bridging¹⁶ has been proposed. However, the quantitative understanding of the interaction, between the hydrate particle roughness and the apparent volume fraction, remains to be explained.

In addition, CP hydrate formation have been observed in previous studies by Sakemoto et al.,¹⁷ Ishida et al.,¹⁸ Karanjkar et al.,¹⁹ Mitarai et al.,²⁰ and recently Martínez de Baños et al.²¹ All previous investigations indicate a growth in two stages: (i) nucleation at the interface and (ii) growth depending upon the temperature, the presence of a surfactant, or properties of the wall surfaces. Our objective is to study a single water drop^{22,23} as a simplified model of an emulsion of water in CP with an oil-soluble surfactant, sorbitan monooleate (Span 80), test the results of Karanjkar et al.,^{14,19} and add new data to understand the complex morphogenesis of this relatively simple system. Few studies are at the scale of a single drop presumably because of the formation of a “halo”, a hydrate crust growing simultaneously on the water drop and also radially on the glass substrate.^{21,24} Here, this difficulty is circumvented using polytetrafluoroethylene (PTFE) surface film to hold the drop. The effect of the influence of inhibitors and anti-agglomerants, such as colloidal particles, polymer surfactants, or salt,¹⁷ has to be taken into account because their interfacial and adhesive properties are known to affect the kinetics of hydrates.^{25,26}

This paper is composed of three parts. In the first part, the fluid used and experimental setup are described together with new measurements of the fluid properties, such as the interfacial tension as a function of the temperature and

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surfactant concentration. In the second part, the growth and dissociation of CP hydrates are observed and monitored for a well-defined temperature protocol. Finally, a summary of the experiments of CP hydrate formation and dissociation times and temperatures for various surfactant concentrations is presented and discussed.

MATERIALS AND METHODS

Fluids. Ultrapure water is drawn from a purification system, which produces water with a constant electric resistivity of 18.2 M Ω cm. The oil phase is CP (reagent grade of 98%). CP is a ring molecule, C₅H₁₀, with a density of 0.751 g L⁻¹, a melting temperature of -93.9 °C (179.25 K), and a boiling temperature of 49.2 °C (322.35 K). Next, a water drop is placed in a bath of CP. At room temperature, CP is known to be weakly soluble in water (156 mg L⁻¹). In most cases, CP is mixed with an oil-soluble surfactant, Span 80, used as received (from Sigma-Aldrich). C₂₄H₄₄O₆ is a non-ionic surfactant with a hydrophilic-lipophilic balance (HLB) of 4 \pm 1. Other sorbitan monooleate surfactants (Span 20, 40, and 60) with different chain lengths have been shown to have an influence on the drop sizes and dissociation temperature.²⁷ In addition, heptane (from VWR Chemicals with a reagent grade of 98%) was used as the test case. The properties of the fluids depend upon the temperature. The densities of CP and water have been measured using a densimeter (DMA 35 from Anton Paar) with an accuracy of 10⁻³ g cm⁻³ and are in good agreement with the previous results for alkanes²⁸ and water.²⁹ The stoichiometric CP/water molar ratio for complete conversion to hydrate is about 1:17 according to Zhang et al.;³⁰ therefore, CP is in excess for the hydrate formation reaction.

Interfacial Tension. Interfacial tension was measured using a drop shape tensiometer (DSA100 manufactured by Krüss). The measure is based on the Young-Laplace model³¹ and the shape of the drop to calculate the interfacial tension, σ . The CP rising drops were formed using an inverted needle (J type) of 2.13 mm outer diameter in water. The water temperature was controlled using a water bath (Julabo F12). Figure 1 presents the interfacial tension measurements for

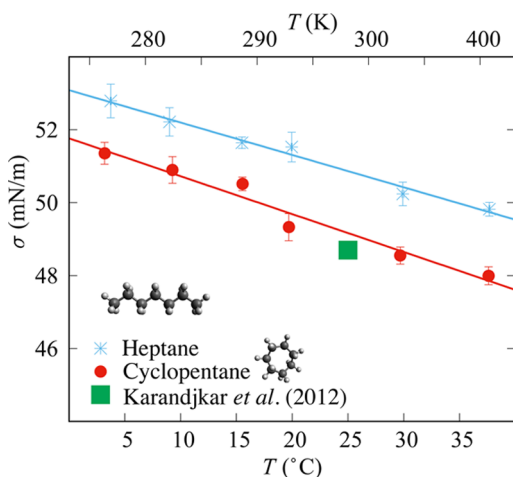


Figure 1. Interfacial tension, σ , versus temperature, T , for heptane and CP. The green square represents the data from Karandjkar et al.¹⁹ The lines are linear fits of the data.

heptane and CP for different temperatures. The results for the interfacial tension of heptane in water are in agreement with Zeppieri et al.³² Errors bars have been added and represent a variance over 15 measurements. For CP, the new data are consistent with the measurement of Karandjkar et al.¹⁹ The interfacial tension of pure fluids is relatively high and decreases with the temperature in a linear fashion.

To quantify the effects of Span 80 on the interfacial tension, mixtures of CP with six different concentrations, C_{S80} , of Span 80 (in

volume) have been prepared. Figure 2 shows the equilibrium interfacial tension of water and CP as a function of C_{S80} over several

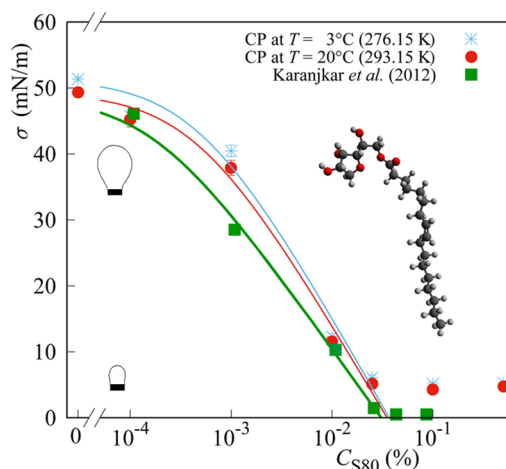


Figure 2. Interfacial tension, σ , versus Span 80 concentration, C_{S80} (v/v), for different temperatures. The lines are Langmuir fits for $C_{S80} < \text{cmc}$. The green squares and line are results by Karandjkar et al.¹⁹ at 25 °C (298.15 K). Insets are drawings of the drop shapes (left) and the molecular structure of Span 80 (right).

decades of concentration. Results are obtained for two temperatures: 3 °C (276.15 K) and 20 °C (293.15 K). The results of Karandjkar et al.¹⁹ are also reproduced in Figure 2 and are lower but in reasonably good agreement the present data. The critical micelle concentration (cmc) can be identified around 0.025% (v/v) (or 6×10^{-4} mol/L), again in good agreement with results of Karandjkar et al.¹⁹ Below cmc, the variation of the interfacial tension, σ , with C_{S80} can be described by the Langmuir adsorption isotherm³³

$$\frac{\Gamma_e}{\Gamma_\infty} = \frac{1}{1 + \left(\frac{\alpha}{\beta C_{S80}}\right)} \quad (1)$$

where Γ_e is the equilibrium instantaneous surface concentration of Span 80, Γ_∞ is the surface concentration of Span 80 (45 \AA^2 per adsorbed Span 80 molecule or $3.7 \times 10^{-6} \text{ mol m}^{-2}$) at saturation, and α/β is the only adjustable parameter representing the kinetic rate constants obtained from least squares methods.

With substitution of Γ_e by Γ from eq 1 in the equation of state for the interfacial tension, the experimental data can be described by

$$\sigma(\Gamma) = \sigma_c + RT\Gamma_\infty \left[\ln \left(1 - \frac{\Gamma}{\Gamma_\infty} \right) \right] \quad (2)$$

where σ_c represents the clean interfacial tension without Span 80, R is the universal gas constant ($8.314 \text{ J mol}^{-1} \text{ K}^{-1}$), and T is the temperature in kelvin.

For large C_{S80} , the present measurements are below the values from Karandjkar et al.¹⁹ This could be explained by the sensitivity of low interfacial tension to tip contact angles. These delicate measurements could be better measured using a spinning drop tensiometer.

Visualization. A Linkam cell (model THMS600) is used that allows for temperature control through a nitrogen gas coolant. It is fixed on a microscope (Olympus BX51). This microscope can be used in transmission or reflection and with white polarized light. Focus 2.5 \times , 10 \times , 20 \times , and 50 \times objectives were used to observe the drops. A schematic of the experimental setup is presented on Figure 3. The diameter of the container is 16 mm, and the height of the container is 2 mm. The water drops were prepared using a micropipette (Eppendorf) with a volume of $2 \pm 0.05 \mu\text{L}$. The drops used here are slightly larger than the drops studied by Ning and Liu²³ and of similar size to those used by Karandjkar et al.¹⁹ However, the nucleation temperature of the water drop in CP is related to the drop diameter,

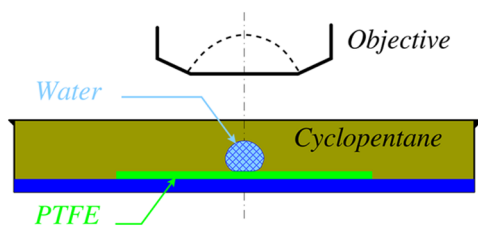


Figure 3. Schematic of the experimental setup. The container is filled with CP, and a single $2 \mu\text{L}$ water drop is formed on the PTFE surface. The temperature protocol is applied, and the observations are from the top.

and the presence of impurities remains variable for every experiment and is in the range from $-15 \text{ }^\circ\text{C}$ (258 K) to $-25 \text{ }^\circ\text{C}$ (248 K). The drop is placed on PTFE disks in a bath of CP. It is pinned by a small laser cut hole that prevents the drop from moving during the temperature protocol. Finally, a coverslip is over the container.

Temperature Protocol. To form hydrates, a temperature protocol^{19,34} is chosen that first converts the water to ice. In the present experiments, the presence of ice is used to trigger the formation of hydrates. The temperature ramping rate is fixed at $5 \text{ }^\circ\text{C}/\text{min}$ and is similar to the ramping rate of previous experiments.^{19,34} The temperature protocol is sketched in Figure 4 and is as follows: (i)

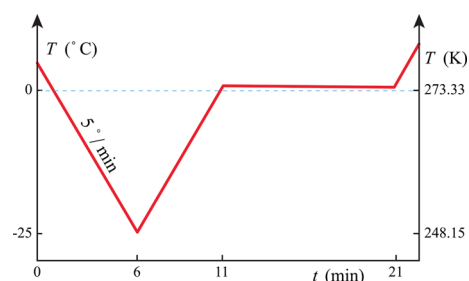


Figure 4. Schematic of the temperature sequence experienced by water drops in CP. Note that the cooling and heating rates are constant at $5 \text{ }^\circ\text{C}/\text{min}$ and the plateau is at $0.2 \text{ }^\circ\text{C}$ (273.15 K).

the temperature, T , decreases from $5 \text{ }^\circ\text{C}$ (273 K) to $-25 \text{ }^\circ\text{C}$ (248 K) at $5 \text{ }^\circ\text{C}/\text{min}$ to convert the water drop into ice; (ii) T increases to $0.2 \text{ }^\circ\text{C}$ (273.35 K), again with a temperature ramp of $5 \text{ }^\circ\text{C}/\text{min}$; (iii) T remains constant at $0.2 \text{ }^\circ\text{C}$ (273.35 K) for 10 min, so that the ice melts and CP hydrates can grow; and (iv) T increases to room temperature, and the system dissociates. Note that the minimum reported CP hydrate equilibrium dissociation temperature varies from 7 to $7.7 \text{ }^\circ\text{C}$ (from 280.15 to 280.85 K) according to Zhang et al.,⁴ Sakemoto et al.,¹⁷ Sloan and Koh,¹ and Aman et al.¹⁶

RESULTS AND DISCUSSION

The results consist of a series of pictures that report the morphology of water drops in CP and a table for the ice nucleation temperature, duration of CP hydrate formation, and dissociation temperature.

CP Hydrates from Ice. Following the temperature sequence described earlier, the water drop reaches three steady states where it (i) freezes into ice balls, (ii) forms hydrates, and (iii) dissociates to liquid drops. Figure 5 presents images of the three steady states. These are bright in the center because of transmitted light going through the center of the drop. The freezing takes place at T between $-17.4 \text{ }^\circ\text{C}$ (255.75 K) and $-19.1 \text{ }^\circ\text{C}$ (254.05 K), and the time scale for complete crystallization is on the order of a few seconds. The critical freezing temperature is stochastic in nature.³⁵ Moreover, the drop is frozen from the bottom; therefore, the propagation of a freezing front and a reduction of mass density explains the pointy ice drop with a sharp tip.³⁶ This tip can be observed in Figure 5a.

During the second temperature ramp, the temperature reaches $0.2 \text{ }^\circ\text{C}$ (273.35 K) and the ice starts to melt. During the temperature plateau at $0.2 \text{ }^\circ\text{C}$ (273.35 K), the ice continues to melt and the hydrates grow. This process takes a few minutes before reaching a steady state. Figure 5b presents an image of the steady state after the ice melted. Experiments³⁷ on larger drops and gas hydrates (methane and carbon dioxide) at elevated pressures have shown long-term hydrate growth (over hours). At least our experiments have reached metastable states, where lateral growth has taken place. Although our experiment only had access to the top view, the ice particle with growing hydrates seems to be expanded. This may be interpreted as the porous shell hydrate formed. The full conversion of the ice into hydrate also depends upon the shell material properties (porosity and stiffness), conversion rate, and diffusion of water through the shell. For water drops in CP, the lateral growth velocity and interface thickness have been investigated before^{14,21,38} and some of the properties of the porous hydrate shell are available. The surface morphology of the hydrate sphere (Figure 5b) differs from the ice ball (Figure 5a), specifically the roughness. In an emulsion, the increase of roughness might lead to an increase in the apparent volume fraction of the disperse phase, hence a significant increase of the viscosity. The surface morphology is known to depend upon the magnitude of the driving force, i.e. the subcooling temperature.^{20,21,37,38} Here, the subcooling temperature is relatively high at about $7 \text{ }^\circ\text{C}$ (280 K). During the last

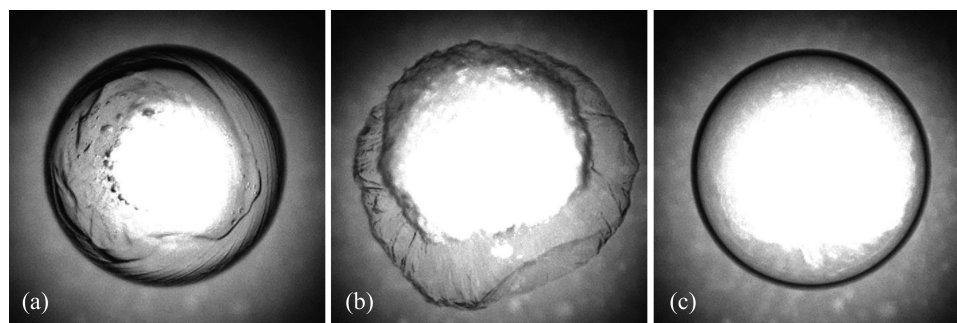


Figure 5. View from the top of the time evolution of a water drop in CP experiencing the temperature sequence from room temperature to (a) $-25 \text{ }^\circ\text{C}$ (248 K), where crystallization occurs, (b) to $0.2 \text{ }^\circ\text{C}$ (273.15 K), where ice melts and hydrate grows, and (c) back to room temperature. The drop volume is $2 \mu\text{L}$, which corresponds to about 1.5 mm in diameter.

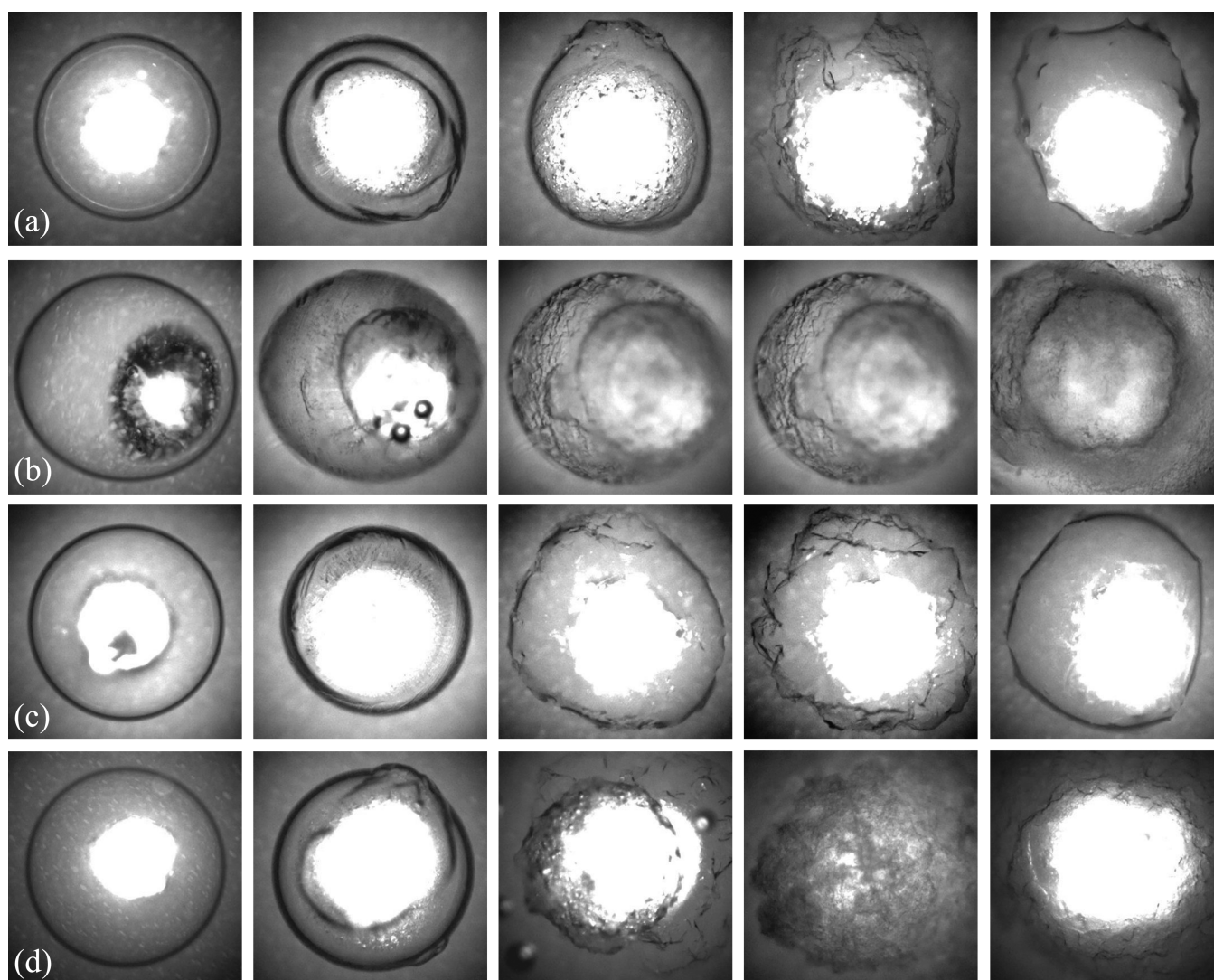


Figure 6. Snapshots of CP hydrate growth for various C_{S80} : (a) 0%, (b) 0.0001%, (c) 0.001%, and (d) 0.1%, the only concentration above cmc. The first image of each line is the initial water drop with an initial volume of $2 \mu\text{L}$, which corresponds to a diameter of about 1.5 mm, at $T = 5 \text{ }^\circ\text{C}$ (278.15 K). The images on the second, third, and fourth positions are at $T = 0.2 \text{ }^\circ\text{C}$ (280 K). The last image of each line illustrates the onset of hydrate melting at a temperature around $7 \text{ }^\circ\text{C}$ (273.35 K).

temperature ramp, the temperature increases and the hydrate melts back to a water drop, as depicted in Figure 5c.

Effect of the Span 80 Surfactant. The temperature sequence has been applied to drops in various baths of CP and CP mixtures with Span 80 of different concentrations, C_{S80} . The morphology of the CP hydrates is tentatively related to the interfacial tension between water and CP, which was measured for different temperatures in Figure 2. Specifically, the time evolution of the drop during the temperature cycle is presented for $C_{S80} = 0, 0.0001, 0.001,$ and 0.1% in panels a, b, c, and d of Figure 6, respectively. Only the last concentration (Figure 6d) is above the cmc. The first image of the line is the initial water drop pinned to the hole of the PTFE disk. Note the slight differences in the droplet area of the initial image, which is due to the effect of the hole used to prevent the drop to move: its shape can vary slightly and modify the position of the contact line and, subsequently, the droplet area. Moreover, the increase in the surfactant concentration is known to flatten the droplet and induce an increase of the droplet area. The second, third, and fourth images render the hydrate growth. The last image of each line illustrates the melting of the hydrate. The lighting of

each line is independent. The thinning and roughening of the hydrate interface is clear, especially for $C_{S80} = 0.0001\%$ (Figure 6b), which corresponds to a concentration below cmc, where the surfactant molecules are not enough to cover the whole drop. For $C_{S80} > \text{cmc}$, presented in Figure 6d, the morphology of the hydrate seems to be slightly more dense, because the light is blocked. However, the detailed roughness morphology of the porous hydrate seems to be similar and independent of C_{S80} . In other words, the low interface tension and associated energy cost of excess interfacial area seems negligible compared to the subcooling driving force, which favors hydrate formation.

During the hydrate growth, conical structures have been observed in the mixture of 0.1% Span 80 in CP. Figure 7 presents images of the growth of conical hydrates also observed in previous experimental studies.^{8,19} These structures can be seen to grow after 16 s through the temperature plateau at $T = 0.2 \text{ }^\circ\text{C}$ (273.35 K). At 22 s, the crystals rise upward, persist, and aggregate to form a bulk hydrate.

Table 1 summarizes the experiments on interfacial tension measurements at different temperature, $\sigma_3 \text{ }^\circ\text{C}$ (275 K) and $\sigma_{20} \text{ }^\circ\text{C}$ (293 K). In addition, the duration of the hydrate formation,

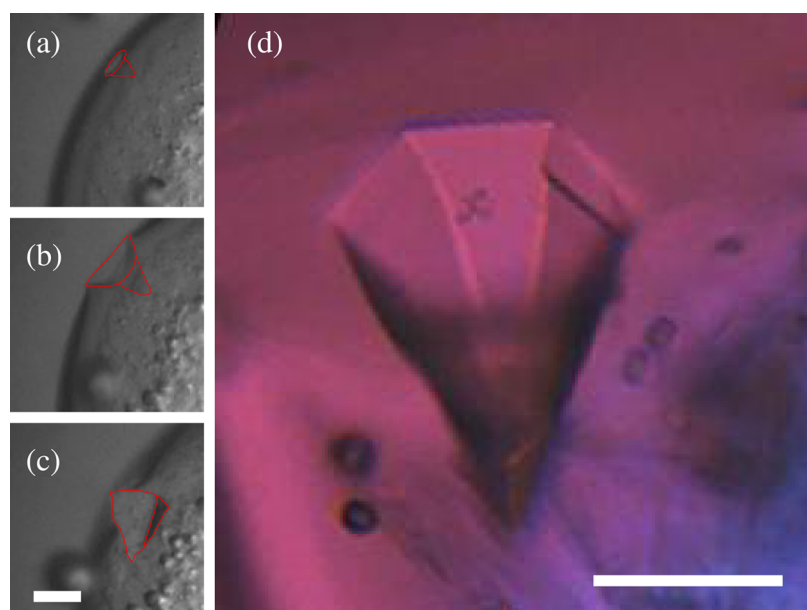


Figure 7. (a–c) Growth of a conical hydrate crystal in a solution of CP with 0.1% Span 80 over 22 s during the temperature plateau at $T = 0.2\text{ }^{\circ}\text{C}$ (273.35 K). The red outlines identify typical hydrate crystal borders, and the scale bar represents $100\text{ }\mu\text{m}$. (d) View of the conical crystal through polarized light microscopy. The scale bar represents $25\text{ }\mu\text{m}$.

Table 1. Summary of the Experiments for Different CP–Span 80 Solutions: Interfacial Tension at Different Temperatures, Time for Hydrate Formation, t_f , at $0.2\text{ }^{\circ}\text{C}$ (273.35 K), Dissociation Temperature, T_d , and Time for Dissociation, t_d

C_{S80} (wt %)	$\sigma_3\text{ }^{\circ}\text{C}$ (275 K) (mN/m)	$\sigma_{20}\text{ }^{\circ}\text{C}$ (293 K) (mN/m)	t_f (s)	T_d [$^{\circ}\text{C}$ (K)]	t_d (s)
0	51.3	49.3	257	6.7 (279.85)	27
0.0001	45.7	45.2	457	5.7 (278.85)	37
0.001	40.4	37.9	467	5.5 (278.65)	39
0.1	5.1	4.3	341	7.7 (280.85)	135

t_f when reaching the plateau temperature at $T = 0.2\text{ }^{\circ}\text{C}$ (273.35 K) is also reported. t_f increases with C_{S80} . This is consistent with the effect of surfactants that will settle at the interface between water and CP and will prevent the motion of water molecules and CP molecules and the formation of hydrate structures. Table 1 also reports information about the dissociation temperature and duration, T_d and t_d . The dissociation temperature remains in a range consistent with previous experiments that found T_d from 7 to $7.7\text{ }^{\circ}\text{C}$ (from 280.15 to 280.85 K) according to Zhang et al.,⁴ Sakemoto et al.,¹⁷ Sloan and Koh,¹ and Aman et al.¹⁶

Moreover, the duration for total dissociation, t_d , increases with C_{S80} . This seems to be directly related to the drop in interfacial tension between water and CP, which was measured for different temperatures in Figure 2.

CONCLUSION

Microscopic observations of the formation and growth of CP hydrate crystals, at atmospheric pressure, were performed through temperature quenches of water drops. These observations were performed systematically using $2\text{ }\mu\text{L}$ water drops in CP baths with a specific temperature protocol that nucleated ice at a low temperature. The simultaneous melting of water and growth of hydrates are observed during a temperature plateau at $0.2\text{ }^{\circ}\text{C}$ (273.35 K) for 10 min.

These morphology data are completed with detailed interfacial tension measurements for different temperatures and surfactant concentrations. A large concentration range of the Span 80 surfactant was tested, and it seems that a slightly rougher porous shell is observed. In addition, the kinetic time scale is increased by the addition of a surfactant. Future studies should be carried out with more realistic systems with several drops, different sizes of drops, and different types of oils and surfactants to make the link between these fundamental experiments and field operations, specifically the adhesion forces and the relation with plugging of pipelines.

AUTHOR INFORMATION

Corresponding Author

*E-mail: jorge.peixinho@univ-lehavre.fr

ORCID

Jorge Peixinho: 0000-0003-0850-3406

Notes

The authors declare no competing financial interest.

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